

# The Hybrid Idea of (Energy Minimization) Optimization Methods Applied to Study Prion Protein Structures Focusing on the $\beta 2$ - $\alpha 2$ Loop

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**Abstract:** In molecular mechanics, current generation potential energy functions provide a reasonably good compromise between accuracy and effectiveness. This paper firstly reviewed several most commonly used classical potential energy functions and their optimization methods used for energy minimization. To minimize a potential energy function, about 95% efforts are spent on the Lennard-Jones potential of van der Waals interactions; we also give a detailed review on some effective computational optimization methods in the Cambridge Cluster Database to solve the problem of Lennard-Jones clusters. From the reviews, we found the hybrid idea of optimization methods is effective, necessary and efficient for solving the potential energy minimization problem and the Lennard-Jones clusters problem. An application to prion protein structures is then done by the hybrid idea. We focus on the  $\beta 2$ - $\alpha 2$  loop of prion protein structures, and we found (i) the species that has the clearly and highly ordered  $\beta 2$ - $\alpha 2$  loop usually owns a  $3_{10}$ -helix in this loop, (ii) a “ $\pi$ -circle” Y128-F175-Y218-Y163-F175-Y169-R164-Y128(-Y162) is around the  $\beta 2$ - $\alpha 2$  loop.

**Key words:** Hybrid idea; computational optimization methods; energy minimization; potential energy; Lennard-Jones clusters; application to protein structures; prion proteins; the  $\beta 2$ - $\alpha 2$  loop.

## 1 Introduction

In molecular mechanics, current potential energy functions provide a reasonably good accuracy to structural data obtained from X-ray crystallography and nuclear magnetic resonance (NMR), and dynamic data obtained from spectroscopy and inelastic neutron scattering and thermodynamic data. Currently, AMBER, CHARMM, GROMOS and OPLS/AMBER are among the most commonly used classical potential energy functions [1, 11, 35, 38, 42]. The energy,  $E$ , is a function of the atomic positions,  $R$ , of all the atoms in the system these are usually expressed in term of Cartesian coordinates. The value of the energy is calculated as a sum of bonded (or internal) terms  $E_{bonded}$ , which describe the bonds, angles and bond rotations in a macromolecule, and a sum of non-

bonded (or external) long-range terms  $E_{non-bonded}$  [1, 11, 42]:

$$\begin{aligned}
E_{potential} &= E_{bonded} + E_{non-bonded} \\
&= (E_{bond-stretch} + E_{angle-bend} + E_{rotate-along-bond} \\
&\quad (+E_{Urey-Bradley} + E_{improper} + U_{CMAP})) \\
&\quad + (E_{van-der-Waals} + E_{electrostatic} + E_{hydrogen-bonds}).
\end{aligned} \tag{1}$$

For example, for AMBER and CHARMM force fields [11, 42] respectively, the potential energy functions are [11]:

$$\begin{aligned}
E_{AMBER} &= \sum_{bonds} k_b(b - b_0)^2 + \sum_{angles} k_\theta(\theta - \theta_{eq})^2 + \sum_{dihedrals} (V_n/2)[1 + \cos(n\phi - \gamma)] \\
&\quad + \sum_{i=1}^{N-1} \sum_{j=i+1}^N \left[ \frac{A_{ij}}{R_{ij}^{12}} - \frac{B_{ij}}{R_{ij}^6} + \frac{q_i q_j}{\epsilon R_{ij}} \right],
\end{aligned} \tag{2}$$

$$\begin{aligned}
E_{CHARMM} &= \sum_{bonds} k_b(b - b_0)^2 + \sum_{angles} k_\theta(\theta - \theta_0)^2 + \sum_{dihedrals} k_\phi[1 + \cos(n\phi - \delta)] \\
&\quad + \sum_{Urey-Bradley} k_u(u - u_0)^2 + \sum_{impropers} k(\omega - \omega_0)^2 + \sum_{\phi, \psi} V_{CMAP}(\phi, \psi) \\
&\quad + \sum_{nonbonded} \left( \epsilon \left[ \left( \frac{R_{min_{ij}}}{R_{ij}} \right)^{12} - \left( \frac{R_{min_{ij}}}{R_{ij}} \right)^6 \right] + \frac{q_i q_j}{\epsilon R_{ij}} \right),
\end{aligned} \tag{3}$$

where  $b, \theta, \phi, R_{ij}, u, \omega, \psi$  are basic variables ( $b$  is the bond length of two atoms,  $\theta$  is the angle of three atoms,  $\phi$  is the dihedral angle of four atoms,  $R_{ij}$  is the distance between atoms  $i$  and  $j$ ), and all other mathematical symbols are constant parameters specified in various force fields respectively. This paper will discuss how to effectively and efficiently use computational optimization methods to solve the minimization problem of the potential energy in Eq. (1), i.e.

$$min \quad E_{potential}. \tag{4}$$

Firstly, for Eq. (4), we consider why we should perform energy minimization (EM). There are a number of reasons:

- (i) To remove nonphysical (or bad) contacts / interactions. For example, when a structure that has been solved via X-ray crystallography, in the X-ray crystallization process, the protein has to be crystallized so that the position of its constituent atoms may be distorted from their natural position and contacts with neighbors in the crystal can cause changes from the in vitro structure; consequently, bond lengths and angles may be distorted and steric clashes between atoms may occur. Missing coordinates obtained from the internal coordinate facility may be far from optimal. Additionally, when two sets of coordinates are merged (e.g., when a protein is put inside a water box) it is possible that there are steric clashes / overlap presented in the resulting coordinate set ([www.charmmtutorial.org](http://www.charmmtutorial.org)).

- (ii) In molecular dynamics (MD) simulations, if a starting configuration is very far from equilibrium, the forces may be excessively large and the MD simulation may fail [1].
- (iii) To remove all kinetic energy from the system and to reduce the thermal noise in the structures and potential energies [1].
- (iv) Re-minimize is needed if the system is under different conditions. For example, in quantum mechanics / molecular mechanics (QM/MM) one part of the system is modeled in QM while the rest is modeled in MD, re-minimize the system with a new condition is needed ([www.charmmtutorial.org](http://www.charmmtutorial.org)).

To perform EM is to make the system reaching to a equilibration state. EM of Eq. (4) can be challenging, as there are many local minima that optimization algorithms might get stuck in without finding the global minima - in most cases, this is what will actually happen. Thus, how much and how far we should minimize should be well considered. Over-minimization can lead to unphysical “freezing” of the structure and move too much from its original conformation; if not minimized enough and exactly, for example, the normal mode calculation cannot arrive at the bottom of its harmonic well. However, in MD, because the output of minimization is to be used for dynamics, it is not necessary for the optimization to be fully converged but a few hundreds or tens of local optimization search are good and kind enough. To make enough local optimization, usually, after we put the protein into a solvent (e.g. waters), first we restrain the protein by holding the solute fixed with strong force and only optimize the solvent, next holding the solute heavy atoms only, and then holding the CA atoms only, and lastly remove all restraints and optimize the whole system.

Secondly, for Eq. (4), we consider what optimization algorithms we should use. In packages of [1, 5, 11] etc, the following three local search optimization methods have been used.

- (i) SD (steepest descent) method is based on the observation that if the real-valued function  $E(x)$  is defined and differentiable in a neighborhood of a point  $x_0$  then  $E(x)$  decreases fastest if one goes from  $x_0$  in the direction of the negative gradient of  $E(x)$  at  $x_0$ . SD method is the simplest algorithm, it simply moves the coordinates in the negative direction of the gradient (hence in the direction of the force - the force is the (negative) derivative of the potential), without consideration of build ups in previous steps - this is the fastest direction making the potential energy decrease. SD is robust and easy to implement. But SD is not the most efficient especially when closer to minimum and in the vicinity of the local minimum. This is to say, SD does not generally converge to a local minimum, but it can rapidly improve the conformation when system is far from a minimum - quickly remove bad contacts and clashes.
- (ii) Conjugate gradient (CG) method is a method adds an orthogonal vector to the current direction of optimization search and then moves them in another direction nearly perpendicular to this vector. CG method is fast-converging and

uses gradient information from previous steps. CG brings you very close to the local minimum, but performs worse far away from the minimum. CG is slower than SD in the early stages of minimization, but becomes more efficient closer to the energy minimum. In GROMACS CG cannot be used with constraints and in this case SD is efficient enough. When the forces are truncated according to the tangent direction, making it impossible to define a Lagrangian, CG method cannot be used to find the EM path.

(iii) L-BFGS method is a Quasi-Newton method that approximates the reverse of Hessian matrix  $[\nabla^2 E(x)]^{-1}$  of  $E(x)$  for the Newton method search direction  $-[\nabla^2 E(x)]^{-1} \nabla E(x)$ . L-BFGS method is mostly comparable to CG method, but in some cases converges 2~3 times faster with super-linear convergent rate (because it requires significantly fewer line search steps than Polak-Ribiere CG). L-BFGS of Nocedal approximates the inverse Hessian by a fixed number of corrections from previous steps. In practice L-BFGS converges faster than CG.

(iv) The combination of CG and LBFGS, so-called lbfgs-TNCG-BFGS method is a preconditioned truncated Newton CG method, it requires fewer minimization steps than Polak-Ribiere CG method and L-BFGS method, but L-BFGS can sometimes be faster in the terms of total CPU times.

If a global optimization is required, approaches such as simulated annealing (SA), parallel tempering method (super SA, also called replica exchange [60]), Metropolis algorithms and other Monte Carlo methods, Simplex method, Nudged Elastic Band method, different deterministic methods of discrete or continuous optimization etc may be utilized. The main idea of SA refinement is to heat up the system such that the molecule of interest has enough energy to explore a wide range of configurational space and get over local optimal energy barriers. Relatively large structural rearrangements are permitted at these high temperatures. As the temperature is cooled gradually, the structural changes proceed in smaller steps, continuing to descend toward the global energy minimum.

For solving Eq. (4), without considering  $E_{hydrogen-bonds} = \sum_{i=1}^{N-1} \sum_{j=i+1}^N \left[ \frac{C_{ij}}{R_{ij}^{12}} - \frac{D_{ij}}{R_{ij}^{10}} \right]$ , about 95% of the CPU time of calculations is spent at

$$\min E_{van-der-Waals} = \sum_{i=1}^{N-1} \sum_{j=i+1}^N \left[ \frac{A_{ij}}{R_{ij}^{12}} - \frac{B_{ij}}{R_{ij}^6} \right], \quad (5)$$

where  $C_{ij}, D_{ij}$  are constants. In [35], this problem is also called Lennard-Jones (LJ) Atomic Cluster Optimization problem (where within the field of atomic clusters only nonbonded interactions are accounted for and particles are considered to be charge-free; e.g. real clusters of metals like gold, silver, and nickel). It is very necessary to up to date review some effective and efficient computational methods for solving Eq. (5). There are numerous algorithms to solve Eq. (5); here we just list the ones in The Cambridge Energy Landscape Database (<http://doye.chem.ox.ac.uk/jon/structures/LJ.html>) which can obtain the best global structures:

- Hoare and Pal’s work [22, 23, 24] may be the early most successful results on LJ problem. The idea is using build-up techniques to construct the initial solutions which are expected to represent low energy states, and using those initial solutions as starting points for a local search method to relax to the optimal solution [23]. The starting seed is the regular unit tetrahedron with atoms at the vertexes, the obvious global optimal solution for  $N = 4$ . Beginning with this tetrahedron, Hoare and Pal (1971, 1972) added one atom at a time to construct a sequence of polytetrahedral structures and at last got good results up to  $N = 66$  [22, 23, 24]. For example, for  $N = 5$  its globally optimal trigonal bi-pyramid (bi-tetrahedron) structure is gotten by adding an atom at the tetrahedral capping position over a triangular face; following the bi-tetrahedron structure, the optimal structure of  $N = 6$  is tri-tetrahedron (another known optimal structure for  $N = 6$  is octahedron (using tetrahedral capping over triangular faces and half-octahedral capping over square faces), which is not a polytetrahedron); for  $N = 7$  its best structure constructed is the pentagonal bi-pyramid, a structure with a five-fold axis of symmetry. Many computer science data structure procedures such as greedy forward growth operator and reverse greedy operator can make the build-up technique work well. The application of methods of studying noncrystalline clusters to the study of “spherical” face centred cubic (fcc) microcrystallites was described in [24]. In [22] the chief geometrical features of the clustering of small numbers of interacting particles were described.
- The data structure of Northby [37] in finding the good starting solution is the lattice based structure. The lattice structures consist of an icosahedral core and particular combinations of surface lattice points. A class of icosahedral packings was by constructed in [36] adding successively larger icosahedral shells in layers around a core central atom; this icosahedral lattice can be described as 20 slightly flattened tetrahedrally shaped fcc units with 12 vertices on a sphere centered at the core atom. Atoms within each triangular face are placed in staggered rows in a two dimensional hexagonal close-packed arrangement. Each atom in the interior of a face in a given shell is a tetrahedral capping position relative to three atoms in the underlying shell. Northby (1987) relaxed the structure of [36] to get his IC and FC multilayer icosahedral lattice structures [37]. The IC lattice can be referred to the FORTRAN code in [54]; it consists of all those sites which will comprise the outer shell of the next complete Mackay [36] icosahedron. FC lattice is a slight modification of IC lattice in that its outer shell maintains icosahedral symmetry and consists of points at the icosahedral vertices and the stacking fault positions of the outer IC shell. Basing on the IC and FC lattices, Northy (1987) gave his algorithm first finding a set of lattice local minimizers and then relaxing those lattice minimizers by performing continuous minimization starting with those lattice minimizers [37]. The algorithm was summarized as Algorithm 1 and Algorithm 2 of [54].
- The great majority of the best known solutions of Northy [37] are icosahedral in character. The hybridization of global search and local search methods, usually, is more effective to solve the large scale problem than the global search method or local search method working alone. Catching those two ideas, Romero et

al. (1999) combined a genetic algorithm with a stochastic search procedure on icosahedrally derived lattices [41, 3]. The structures of the optimal solutions gotten in [41] are either icosahedral or decahedral in character. The best results of [50] for  $N = 82, 84, 86, 88, 92, 93, 94, 95, 96, 99, 100$  were gotten by using a genetic algorithm alone. Deaven et al. (1996) also using the genetic algorithm got the optimal value known for the magic number  $N = 88$  [16].

- The successful works to improve Northby’s results in [37] were mainly done by Xue [53, 54], Leary [33], and Doye et al. [18, 17].
  - Xue (1994a) introduced a modified version of the Northby algorithm [53]. He showed that in some cases the relaxation of the outer shell lattice local minimizer with a worse potential function value may lead to a local minimizer with a better value. In Northby’s algorithm [37] the lattice search part is a discrete optimization local search procedure, which makes a lattice move to its neighboring lattice with  $O(N^{\frac{2}{3}})$  time complexity. In [53] Xue (1994a) introduced a simple storage data structure to reduce the time complexity to  $O(N^{\frac{2}{3}})$  per move; and then used a two-level simulated annealing algorithm within the supercomputer CM-5 to be able to solve fastly the LJ problem with sizes as large as 100,000 atoms. In [54] by employing AVL trees [25] data structure Xue (1994b) furthermore reduced the time complexity to  $O(\log N)$  if NN (nearest neighbor) potential function is used. Xue (1994b) relaxed every lattice local minimizer found instead of relaxing only those lattice local minimizers with best known potential function value by a powerful Truncated Newton local search method [54], and at last got the best results known for  $N = 65, 66, 134, 200, 300$ .
  - Leary (1997) gave a successful Big Bang Algorithm [33] for getting the best values known of  $N = 69, 78, 88, 107, 113, 115$ . In [33], the FCC lattice structure is discussed and its connections are made with the macrocluster problem. It is also concluded in [33] that almost all known exceptions to global optimality of the well-known Northby multilayer icosahedral conformations for microclusters are shown to be minor variants of that geometry. The Big Bang Algorithm contains 3 steps: Step 1 is an initial solution generating procedure which randomly generates each coordinate of the initial solution with the independently normal distribution; Step 2 is to generate the new neighborhood solution by discrete-typed fixed step steepest descent method, which is repeated until no further progress is made; Step 3 is to relax the best solution gotten in Step 2 by a continuous optimization method—conjugate gradient method.
  - Doye et al. (1995) investigated the structures of clusters by mapping the structure of the global minimum as a function of both cluster size and the range of the pair potential which is appropriate to the clusters of diatomic molecule,  $C_{60}$  molecule, and the ones between them both [17]. For the larger clusters the structure of the global minimum changes from icosahedral to decahedral to fcc as the range is decreased [17]. In [18], Doye et al. (1995) predicted the growth sequences for small decahedral and fcc clusters by

maximisation of the number of NN contacts.

- Calvo et al. (2001) gave some results on quantum LJ Clusters in the use of Monte Carlo methods [9].
- Xiang et al. (2004a) presented an efficient method based on lattice construction and the genetic algorithm and got global minima for  $N = 310\sim 561$  [52] In 2004, Xiang et al. (2004b) continued to present global minima for  $N = 562\sim 1000$  [51].
- Barron-Romero (2005) found the best solutions for  $N = 542-3, 546-8$  in the use of a modified peeling greedy search method [4].
- Takeuchi (2006) found best solutions for  $N = 506, 521, 537-8$  and 541 by a clever and efficient method “using two operators: one modifies a cluster configuration by moving atoms to the most stable positions on the surface of a cluster and the other gives a perturbation on a cluster configuration by moving atoms near the center of mass of a cluster” [45].
- Lai et al. (2011a) found best solutions for  $N = 533$  and 536 using the dynamic lattice searching method with two-phase local search and interior operation [31, 32, 55].
- Algorithms to get the structures at the magic numbers  $N = 17, 23, 24, 72, 88$  (the exceptions to [41]):
  - Freeman et al. (1985) presented the best value for  $N = 17$  when the thermodynamic properties of argon clusters were studied by a combination of classical and quantum Monte Carlo methods [20]. The poly-icosahedral growth of Farges et al. (1985) starts from a 13-atom primitive icosahedron containing a central atom and 12 surface atoms [19]. On each one of the five tetrahedral sites, surrounding a particular vertex, a new atom is added and finally a sixth atom is placed on top to create a pentagonal cap. In this way a 19-atom structure being made of double interpenetrating icosahedra, which is a 13-atom icosahedra sharing 9 atoms, is obtained; i.e., for three pentagonal bipyramids each one shares an apex with its nearest neighbour. In this way a 23-atom model consisting of three interpenetrating icosahedra is gotten for the best value known.
  - Wille (1987) used the SA method yielding low-lying energy states whose distribution depends on the cooling rate to find the best solution known for  $N = 24$  [49].
  - Coleman et al. (1997) proposed a build-up process to construct the optimal solution structures. The HOC (half octahedral cap) structure of the optimal solution for  $N = 72$  is found by a prototype algorithm designed using the anisotropic effective energy simulated annealing method at each build-up stage ([14]).
  - Wales & Doye (1997) gave the lowest values known for  $N = 192, 201$  [46]. Their method is so-called basin-hopping method, in which first the transformed function  $\tilde{f}(x) = \min\{f(x)\}$  was defined and performed starting from

$x$  by the PR conjugate gradient method and then the energy landscape for the function  $\tilde{f}(x)$  was explored using a canonical Monte Carlo simulation.

- Leary (2000) has developed techniques for moving along sequences of local minima with decreasing energies to arrive at good candidates for global optima and got the best value known on  $N = 185$ .

Now we have the outline of some successful optimization methods used to solve Eq.s (4)~(5). We have found the hybrid idea of optimization methods was not emphasized very much (especially for solving Eq. (5)). Thus, in Section 2 of this paper we will emphasize the hybrid idea of optimization methods by introducing our own hybrid methods used to solve Eq.s (4)~(5). Section 3 will present our recent results of applying the hybrid idea of SD and CG and SD again to do EM of some NMR and X-ray prion protein structures in the PDB Bank ([www.rcsb.org](http://www.rcsb.org)); interesting findings will be reported in this Section. Why we choose prion proteins in this study is due to prions effect humans and almost all animals for a major public health concern (e.g. milks and meats we daily drink and eat). At last, in Section 4, we give a concluding remark on the effective and efficient hybrid idea of optimization methods.

## 2 The hybrid idea and some hybrid optimization methods

In this Section, we use how to construct molecular structures of prion amyloid fibrils at AGAAAAGA segment as an example to illuminate the hybrid idea and some hybrid optimization methods we designed.

Neurodegenerative amyloid diseases such as Alzheimer’s, Parkinson’s and Huntington’s all featured amyloid fibrils. Prions also cause a number of neurodegenerative diseases too. All these amyloid fibrils in 3-dimensional quaternary structure have 8 classes of steric zippers, with strong van der Waals interactions between  $\beta$ -sheets and hydrogen bonds between  $\beta$ -strands. Currently, there is no structural information about prion AGAAAAGA amyloid fibrils because of unstable, noncrystalline and insoluble nature of this region, though numerous laboratory experimental results have confirmed this region owning an amyloid fibril forming property (initially described in 1992 by Stanley B. Prusiner’s group). We also did accurate computer calculations on this region and confirmed the amyloid fibril property in this palindrome region [67, 73].

In [58], we constructed three models, model 1 belongs to Class 7 (antiparallel, face=back, up-up) and models 2–3 belong to Class 1 (parallel, face-to-face, up-up) of steric zippers. The models were firstly optimized by SD and then followed by CG. SD has fast convergence but it is slow when close to minimums. CG is efficient but its gradient RMS and GMAX gradient do not have a nice convergence. When the models could not be optimized furthermore, we employed standard SA method (that simulates



the annealing process of crystal materials with Monte Carlo property). After SA, we refined the models by SD and then CG again. SA is a global search optimization method [64] that can make local optimal jump out of / escape from the local trap. We found the refinement results in a loss of potential energy nearly the same magnitude as that of SA; this implies to us SA is very necessary and very effective in our molecular modeling. Numerical results show to us the hybrid is very necessary, effective and efficient.

When the gradient or its generalizations of the target/objective function  $E(x)$  are very complex in form or they are not known, derivative-free methods benefit optimization problems. In [70], we introduced derivative-free discrete gradient (DG) method [2] into the derivative-free global search SA optimization method or genetic algorithms (GAs, which simulate the process natural competitive selection, crossover, and mutation of species), and designed hybrid methods SADG, GADG. In implementation, the hybrids of DG + SADG / GADG + DG were used, and at last SD+CG + SA + SD+CG of Amber package [11] were used to refine the models. We found the hybrids work very well, and more precise best solutions for  $N = 39, 40, 42, 48, 55, 75, 76$ , and 97 were found and their figures show that their structures are more stable than the ones currently best solutions known. We also found the hybrid of evolutionary computations with simulate annealing SA-SAES( $\mu + \lambda$ ), SA-SACEP perform better than evolutionary computations or SA work alone [59].

Canonical dual theory in some sense is the hybrid of the primal and the dual. In [68], we solved the dual problem and then got the solutions for the primal problem. We found the refinement using AMBER package is not necessary. This implies to us the hybrid of primal and dual in canonical dual theory is good enough and effective.

As said in Section 1, in some cases, CG cannot be used to find the EM path; this point will also be shown in next Section (see Tab. 1). Thus, in [67], we specially studied and implement the LBFGS method designed by us and then hybridize it with the LBFGS method of AMBER package. We found the hybrid is very necessary and effective.

By our numerical experiences shown above, the hybrid idea is very necessary, effective and efficient for some hybrid optimization methods in known packages or designed by us. In next Section, we will apply the hybrid idea to do some practical works for some important prion protein NMR and X-ray structures deposited in the PDB Bank.

### 3 An application to prion protein structures, focusing on the $\beta 2$ - $\alpha 2$ loop

Before we use the structure taken from PDB Bank, usually we need to relax it, in order to remove bad contacts and also fix up hydrogen positions. Fairly short local optimization is sufficient to refine and relax the structure. We will use SD-CG-SD local optimization methods. In SD, its search direction is a n-dimensional search and its step-length search is a 1-dimensional search. In CG, the search is usually in a 2-dimensional subspace and conjugacy is a good property only associated with exact line search [43]. Using the hybrid of SD and CG is also in order to remove all these (dimensional) unbalances. In our EM here, the free package Swiss-PdbViewer 4.1.0 (spdbv.vital-it.ch) that has been developed for 20 years is used, we set 3000 steps for SD, then 3000 steps for CG, and then 3000 steps for SD again, Bonds, Angles, Torsions, Improper, Non-bonded and Electrostatic are considered, 12.000 Å is chosen for the Cutoff, stop SD or CG when delta E between two steps is below 0.005 kJ/mol, and stop SD or CG when Force acting on any atom is below 1.000.

We found, for the research of prion proteins, the S2-H2 loop (and its interactions with the C-terminal of H3) is a focus [6, 7, 8, 10, 12, 13, 15, 21, 27, 34, 47, 48, 28, 39, 40, 44, 56, 71, 29, 30, 26]. All prion protein structures have high similarity in three  $\alpha$ -helices (H1, H2, H3) and two  $\beta$ -strands (S1, S2), but there is a great difference just at this S2-H2 loop:

(i) structure with disordered S2-H2 loop:

- mousePrP (1AG2.pdb at 25 °C),
- humanPrP (1QLX.pdb),
- bovinePrP (1DWY.pdb),
- SyrianHamsterPrP (1B10.pdb),
- dogPrP (1XYK.pdb) (- resist to prion infection),
- catPrP (1XYJ.pdb),
- sheepPrP (1UW3.pdb),
- mousePrP[N174T] (1Y15.pdb),
- humanPrP[Q212P]-M129 (2KUN.pdb),
- humanPrP-M129 (1QM1.pdb),
- rabbitPrP[S173N] (2JOH.pdb),
- rabbitPrP[I214V] (2JOM.pdb),
- rabbitPrP[S170N] (4HLS.pdb),
- rabbitPrP[S174N] (4HMM.pdb),
- rabbitPrP[S170N,S174N] (4HMR.pdb),

(ii) structure with highly and clearly ordered S2-H2 loop:

- mousePrP (2L39.pdb at 37 °C),
- mousePrP[V166A] (2KFO.pdb),
- mousePrP[D167S] (2KU5.pdb at 20 °C),
- mousePrP[D167S,N173K] (2KU6.pdb),
- mousePrP[Y169G] (2L1D.pdb),
- mousePrP[Y169A] (2L40.pdb),
- mousePrP[S170N] (2K5O.pdb),
- mousePrP[S170N,N174T] (1Y16.pdb),
- mousePrP[F175A] (2L1E.pdb),
- mousePrP[Y225A,Y226A] (2KFM.pdb),
- mousePrP[Y169A,Y225A,Y226A] (2L1K.pdb at 20 °C),
- elkPrP (1XYW.pdb),
- pigPrP (1XYQ.pdb),
- BankVolePrP (2K56.pdb),
- TammarWallabyPrP (2KFL.pdb),
- rabbitPrP (2FJ3.pdb, 3O79.pdb),
- horsePrP (2KU4.pdb),

where elk and Bank Vole can be infected by prions though they have a highly and clearly ordered S2-H2 loop, and the codes in the brackets are the PDB codes in the PDB Bank. For all these NMR and X-ray structures we did SD-CG-SD relaxation and the variations of the EMs are listed in Tab. 1. From Tab. 1, we can see the energy decreases from SD to CG and from CG to SD. For mousePrP[Y169A], CG is not working well, but it adjusts the SD methods so that it make SD-CG-SD work very well in the second round.

After the SD-CG-SD relaxation of all the structures, now these optimized structures can be used to obtain some helpful structural information [e.g. (i) hydrogen bonds (see Tab. 2), (ii) electrostatic charge distributions on the protein structure surface (see Fig. 1 and Tab. 4), (iii) salt bridges (see Tab. 3), and (iv)  $\pi$ - $\pi$ -stacking and  $\pi$ -cations (see Tab. 5); here why we consider the information of (i)~(iv) is due to “the performance of protein biological function is driven by a number of non-covalent interactions such as hydrogen bonding, ionic interactions, Van der Waals forces, and hydrophobic packing” ([en.wikipedia.org/wiki/Protein\\_structure](http://en.wikipedia.org/wiki/Protein_structure))] at the S2-H2 loop, in order to furthermore understand the S2-H2 loop: why some species has a clearly and highly ordered S2-H2 loop and why some species just has a disordered S2-H2 loop. (i) Using the VMD package ([www.ks.uiuc.edu/Research/vmd/](http://www.ks.uiuc.edu/Research/vmd/)), with Tab. 2, we may observed that the species owning the disordered S2-H2 loop usually does not have a  $3_{10}$ -helix in the S2-H2 loop (except for sheepPrP, rabbitPrP[S170N]-X-ray, rabbitPrP[S174N]-X-ray, and rabbitPrP[S170N,S174N]-X-ray), but the species that has the clearly and highly ordered S2-H2 loop usually owns a  $3_{10}$ -helix, constructed by the following hydrogen bond(s) respectively:

Table 1: Energy variations during the energy minimizations (number of iterations are in the brackets):

Species	SD	CG	SD
mousePrP	-7234.507 (391)	-7243.597 (13)	-7275.160 (45)
humanPrP	-7460.885 (296)	-7610.384 (195)	-7640.157 (54)
bovinePrP	-7698.001 (544)	-7809.687 (255)	-7819.315 (26)
SyrianHamsterPrP	-7418.702 (225)	-7653.251 (258)	-7688.044 (50)
dogPrP	-7148.517 (549)	-7225.084 (151)	-7251.191 (60)
catPrP	-6935.915 (192)	-7186.961 (186)	-7382.361 (229)
sheepPrP	-8066.183 (291)	-8066.300 (22)	-8204.010 (179)
mousePrP[N174T]	-7418.443 (211)	-7657.886 (206)	-7923.308 (419)
humanPrP[Q212P]-M129	-7662.032 (464)	-7688.836 (49)	-7846.123 (424)
humanPrP-M129	-7249.798 (350)	-7345.074 (114)	-7410.470 (99)
rabbitPrP[S173N]-NMR	-7173.492 (271)	-7666.456 (492)	-7684.357 (36)
rabbitPrP[I214V]-NMR	-7785.640 (774)	-7802.710 (39)	-7835.354 (88)
rabbitPrP[S170N]-X-ray	-8682.104 (414)	-8753.827 (178)	-8827.220 (254)
rabbitPrP[S174N]-X-ray	-8551.921 (286)	-8659.515 (160)	-8739.475 (178)
rabbitPrP[S170N,S174N]-X-ray	-8864.615 (363)	-8899.511 (65)	-8945.262 (106)
mousePrP - at 37 °C	-7679.395 (212)	-8031.103 (397)	-8134.173 (213)
mousePrP[V166A]	-8040.172 (436)	-8153.529 (174)	-8164.996 (22)
mousePrP[D167S] - at 20 °C	-7938.545 (594)	-7987.788 (83)	-8051.072 (121)
mousePrP[D167S,N173K]	-7615.915 (546)	-7751.804 (205)	-7906.153 (290)
mousePrP[Y169G]	-7713.983 (249)	-7885.343 (147)	-7913.277 (27)
mousePrP[Y169A]	-7948.267 (507)	-7949.964 (1)	-7955.516 (3)
	-7972.101 (17)	-8064.175 (142)	-8143.807 (179)
mousePrP[S170N]	-7341.668 (74)	-7790.935 (262)	-7947.893 (224)
mousePrP[S170N,N174T]	-7988.545 (468)	-8126.978 (241)	-8149.948 (55)
mousePrP[F175A]	-7660.438 (680)	-7859.632 (381)	-7873.355 (34)
mousePrP[Y225A,Y226A]	-7457.356 (244)	-7588.338 (109)	-7757.797 (213)
mousePrP[Y169A,Y225A,Y226A] - at 20 °C	-7609.773 (173)	-7690.231 (45)	-7700.162 (4)
elkPrP	-7894.305 (875)	-7959.371 (160)	-7978.686 (53)
pigPrP	-6354.886 (141)	-6735.155 (321)	-6813.048 (109)
bankVolePrP	-7727.260 (478)	-7799.991 (118)	-7951.502 (376)
tammarWallabyPrP	-8028.082 (393)	-8195.817 (248)	-8238.248 (69)
rabbitPrP-NMR	-7712.972 (814)	-7730.035 (36)	-7790.620 (154)
rabbitPrP-X-ray	-8643.309 (602)	-8658.327 (40)	-8717.344 (187)
horsePrP	-7335.273 (226)	-7614.526 (433)	-7633.569 (45)

- V166–Y169 - mousePrP at 37 °C, mousePrP[D167S,N173K], mousePrP[F175A], rabbitPrP-X-ray,
- R164–Q168 - mousePrP[V166A],
- P165–Q168 - mousePrP[Y225A,Y226A], elkPrP, BankVolePrP,
- P165–Q168, I166–Y169 - TammarWallabyPrP.

(ii) Seeing Fig. 1 and Tab. 4, we may know that at the S2-H2 loop it is mainly covered by the electrical cloud of negatively (in red color) charged residues [except for mousePrP[D167S,N173K], rabbitPrP-NMR and horsePrP etc (Fig. 2)], with positively (in blue) charged residues R164 (for all species) and H177 (for mousePrP, humanPrP-M129, mousePrP[D167S,N173K] only) at the N-terminal end and C-terminal end of S2-H2 loop respectively (we found there is a salt bridge R164–D178 linking this loop of rabbitPrP-NMR, rabbitPrP-X-ray, horsePrP, dogPrP, elkPrP and buffaloPrP for long time MD simulations [57, 61, 62, 69, 74, 72, 66, 63]). From Tab. 4, we might see that the negatively charged S2-H2 loop might have long distance nuclear overhauser effect (NOE) interactions with the positively charged residues such as K204, R208, K/R220, R227, R228, R229, and R230 at the C-terminal end of H3. (iii) The salt bridges in Tab. 3 might be not very strong and will be quickly broken in a long time MD simulations [62]. (iv) Lastly, we present some bioinformatics of  $\pi$ - $\pi$ -stacking and  $\pi$ -cations (one kind of van der Waals interactions) at the S2-H2 loop. Seeing Tab. 5, we may know at

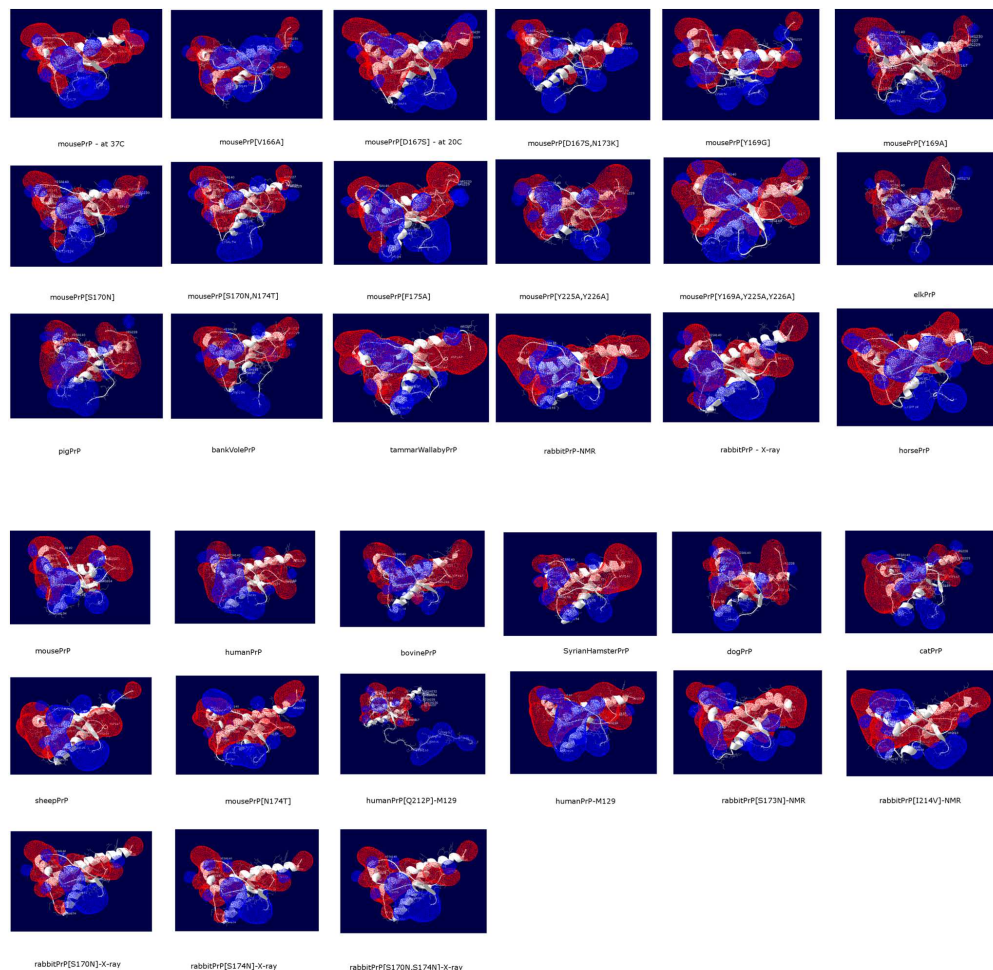


Figure 1: Positively (in blue) and negatively (in red) charged residues distributed on each protein structure surface of the 33 PrPs. The first three rows of PrPs are with highly and clearly ordered S2–H2 loop; but the last three rows of PrPs are with disordered S2–H2 loop.

S2-H2 loop and its contacts with the C-terminal end of H3 there are the following  $\pi$ - $\pi$ -stacks Y169–F175, F175–Y218, Y163–Y218, and the following  $\pi$ -cations R164–Y169, R164–Y128, which clearly contribute to the clearly and highly ordered S2-H2 loop structures [65]. For buffaloPrP, we found another two  $\pi$ -stackings: Y163–F175–Y128 [65, 66]. Thus, for PrPs, we found an interesting “ $\pi$ -circle” Y128–F175–Y218–Y163–F175–Y169–R164–Y128(–Y162) around the S2-H2 loop.

## 4 A concluding remark

In optimization, especially for solving large scale or complex or both optimization problems, the hybrid of optimization (local search or global search) methods is very necessary, and very effective and efficient for solving optimization problems. In molecular mechanics, to optimize its potential energy, even just one part of it e.g. the

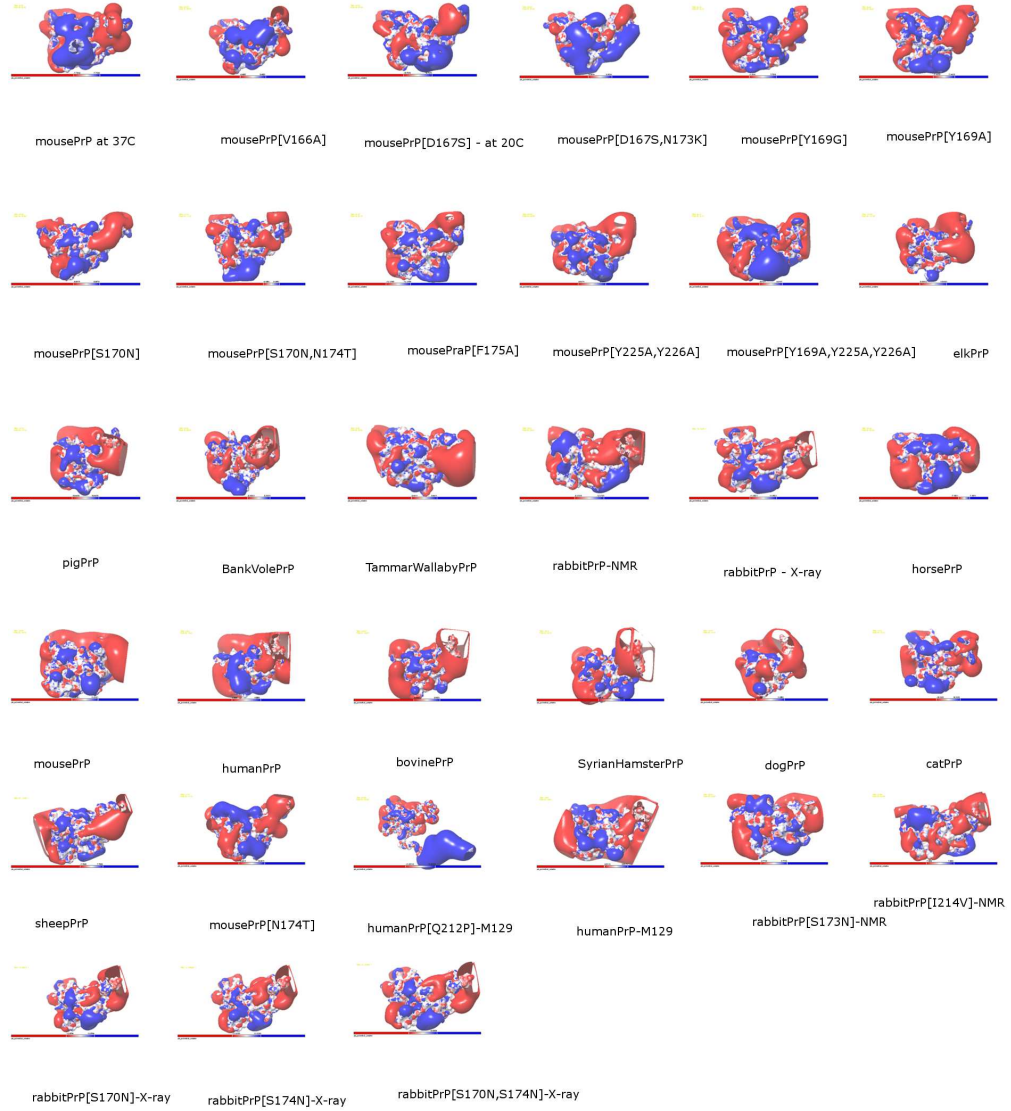


Figure 2: Surface electrostatic charge distributions for each of the 33 PrPs. The first three rows of PrPs are with highly and clearly ordered S2–H2 loop; but the last three rows of PrPs are with disordered S2–H2 loop. Blue is for positive charge whereas red is for negative charge.

Lennard-Jones potential, is still a challenge to optimization methods; the hybrid idea is very helpful and useful. An application to prion protein structures is then done by the hybrid idea. Focusing on the  $\beta 2$ – $\alpha 2$  loop of prion protein structures, we found (i) the species that has the clearly and highly ordered  $\beta 2$ – $\alpha 2$  loop usually owns a  $3_{10}$ -helix in this loop, (ii) a “ $\pi$ -circle” Y128–F175–Y218–Y163–F175–Y169–R164–Y128(–Y162) is around the  $\beta 2$ – $\alpha 2$  loop. In conclusion, this paper proposes a hybrid idea of optimization methods to efficiently solve the potential energy minimization problem and the LJ clusters problem. We first reviewed several most commonly used classical potential

energy functions and their optimization methods used for energy minimization, as well as some effective computational optimization methods used to solve the problem of Lennard-Jones clusters. In addition, we applied this hybrid idea to construct molecular structures of prion amyloid fibrils at AGAAAAGA segment, by which we provided the additional insight for the  $\beta$ 2- $\alpha$ 2 loop of prion protein structures. This study should be of interest to the protein structure field.

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Table 2: Hydrogen bonds at the S2-H2 loop, and its linkage with the C-terminal end of H3 (in the brackets are the distances of the hydrogen bonds):

Species	at the S2-H2 loop	linking with the C-terminal end of H3
mousePrP	D178.OD1-R164.2HH1 (1.97 Å)	S170.OG-Y225:HH (1.71 Å)
humanPrP		
bovinePrP		
SyrianHamsterPrP	Y169.HH-D178.OD2 (1.80 Å) T183.OG1-Y162.H (1.74 Å)	
dogPrP	D167.O-S170.H (1.82 Å)	
catPrP		
sheepPrP	P165.O-Q168.H (1.90 Å) D178.OD2-Y128.HH (1.81 Å) Y163.O-M129.H (1.79 Å) N171.O-N171.2HD2 (1.84 Å) T183.OG1-Y162.H (1.92 Å)	Y163.OH-Q217.2HE2 (1.91 Å)
mousePrP[N174T]	D178.OD2-Y128.HH (1.72 Å) Y163.H-M129.O (1.79 Å)	
humanPrP[Q212P]-M129	N171.OD1-N173.H (1.84 Å) Y162.H-T183.OG1 (1.94 Å)	Q172.OE1-Y218.HH (1.68 Å)
humanPrP-M129	Y162.H-T183.OG1 (1.89 Å)	
rabbitPrP[S173N]-NMR		E220.OE1-Y162.HH (1.68 Å) D177.OD2-Y127.HH (1.71 Å)
rabbitPrP[I214V]-NMR	Y168.O-N170.2HD2 (1.86 Å) Q171.O-V175.H (1.70 Å)	
rabbitPrP[S170N]-X-ray	P165.O-Q168.H (1.93 Å) V166.O-Y169.H (1.92 Å)	Q172.OE1-Q219.2HE2 (1.85 Å)
rabbitPrP[S174N]-X-ray	P165.O-Q168.H (1.94 Å) V166.O-Y169.H (1.92 Å) Y169.HH-D178.OD2 (1.74 Å) N171.HD2-N174.ND2 (1.95 Å)	
rabbitPrP[S170N,S174N]-X-ray	P165.O-Q168.H (1.92 Å) V166.O-Y169.H (1.93 Å) Y169.HH-D178.OD2 (1.78 Å) N171.OD1-N174.H (1.88 Å)	Q172.OE1-Q219.2HE2 (1.89 Å)
mousePrP - at 37 °C	Y163.O-M129.H (1.71 Å) R164.O-Y169.HH (1.80 Å) V166.O-Y169.H (1.83 Å) D178.OD2-Y128.HH (1.74 Å)	S170.O-Y218.HH (1.61 Å)
mousePrP[V166A]	R164.HE-Q168.OE1 (1.90 Å) T183.OG1-Y162.H (2.02 Å)	
mousePrP[D167S] - at 20 °C	R164.2HH2-Q168.OE1 (1.90 Å) V166.O-Y169.H (1.87 Å) T183.OG1-Y162.H (1.90 Å)	
mousePrP[D167S,N173K]	R164.CZ-Q168.1HE2 (1.96 Å) V166.O-Y169.H (1.82 Å) V166.O-S170.H (1.75 Å)	
mousePrP[Y169G]	T183.OG1-Y162.H (1.95 Å)	
mousePrP[Y169A]	P165.O-S170.HG (1.91 Å) T183.OG1-Y162.H (1.94 Å)	Q168.O-Y225.HH (1.73 Å)
mousePrP[S170N]	N170.2HD2-N171.OD1 (1.91 Å) T183.OG1-Y162.H (1.84 Å)	
mousePrP[S170N,N174T]	Y169.HH-D178.OD2 (2.01 Å) N171.O-N171.2HD2 (1.81 Å) T183.OG1-Y162.H (1.79 Å)	
mousePrP[F175A]	V166.O-Y169.H (1.76 Å) N171.1HD2-N174.ND2 (1.75 Å) N171.2HD2-N171.O (1.78 Å)	R164.O-Y218.HH (1.61 Å)
mousePrP[Y225A,Y226A]	D178.OD2-Y128.HH (1.61 Å) Y163.O-MET129.H (1.77 Å) P165.O-Q168.2HE2 (1.85 Å) S170.OG-N174.1HD2 (1.95 Å) N171.H-N174.OD1 (1.94 Å) N173.2HD2-HIS177.NE2 (1.93 Å) T183.OG1-Y162.H (1.96 Å)	
mousePrP[Y169A,Y225A,Y226A] - at 20 °C		
elkPrP	P165-Q168.H (1.80 Å) Y169.HH-D178.OD2 (1.73 Å) T183.OG1-Y162.H (1.79 Å)	
pigPrP		
BankVolePrP	P165.O-Q168.H (1.99 Å) T183.OG1-Y162.H (1.92 Å)	
TammarWallabyPrP	P165.O-Q168.N (1.98 Å) I166.O-Y169.H (1.92 Å) T183.OG1-M162.H (1.79 Å)	
rabbitPrP-NMR	T182.OG1-Y161.H (1.81 Å)	
rabbitPrP-X-ray	V166.O-Y169.H (1.95 Å) D178.OD2-Y169.HH (1.82 Å)	
horsePrP	D178.OD2-Y169.HH (1.71 Å)	

Table 3: Salt bridges at the S2-H2 loop, and its linkage with the C-terminal end of H3 (in the brackets are the distances of the salt bridges):

Species	at the S2-H2 loop	linking to the C-terminal end of H3
mousePrP	D178.OD1-R164.NE (4.01 Å) D178.OD1-R164.NH1 (2.96 Å)	
humanPrP		
bovinePrP		
SyrianHamsterPrP		
dogPrP		
catPrP	D178.OD1-R164.NH2 (4.59 Å) D178.OD2-R164.NE (2.97 Å) D178.OD2-R164.NH1 (4.76 Å) D178.OD2-R164.NH2 (3.06 Å)	
sheepPrP		
mousePrP[N174T]		
humanPrP[Q212P]-M129	E168.OE1-R164.NH1 (3.07 Å) E168.OE1-R164.NH2 (3.33 Å)	
humanPrP-M129	D178.OD2-R164.NH1 (2.99 Å) D178.OD2-R164.NH2 (3.13 Å) D178.OD2-R164.NE (4.75 Å) D178.OD1-R164.NH1 (4.61 Å)	
rabbitPrP[S173N]-NMR		
rabbitPrP[I214V]-NMR		
rabbitPrP[S170N]-X-ray	D178.OD2-R164.NH2 (3.16 Å) D178.OD2-R164.NE (4.14 Å) D178.OD1-R164.NH2 (4.06 Å) D178.OD1-H177.ND1 (2.95 Å)	
rabbitPrP[S174N]-X-ray		
rabbitPrP[S170N,S174N]-X-ray		
mousePrP - at 37 °C		
mousePrP[V166A]		
mousePrP[D167S] - at 20 °C		
mousePrP[D167S,N173K]		
mousePrP[Y169G]		
mousePrP[Y169A]		
mousePrP[S170N]		D167.OD1-R229.NH1 (3.03 Å) D167.OD1-R229.NH2 (3.81 Å)
mousePrP[S170N,N174T]		
mousePrP[F175A]		
mousePrP[Y225A,Y226A]		
mousePrP[Y169A,Y225A,Y226A] - at 20 °C		
elkPrP		
pigPrP	D178.OD2-R164.NH1 (3.00 Å)	
BankVolePrP		
TammarWallabyPrP		
rabbitPrP-NMR		H176.ND1-E210.OE1 (2.96 Å)
rabbitPrP-X-ray		
horsePrP		

Table 4: Cliques of positively charged residues distributed on the surface of each optimized structure of the 33 PrPs (one bracket is one clique):

mousePrP	(R164), (H177), (K220), (R156), (R148), (K185), (R151, R136, K194)
humanPrP	(R164), (R228), (R220), (H155, R156, K185, H187, K194,)
bovinePrP	(R164), (K185), (K194), (H155, R156, R136, H187, R220)
SyrianHamsterPrP	(R164), (K185), (K194), (K220), (R136, R156), (R148, R151)
dogPrP	(R164), (R228), (R220), (K185), (K194), (R151, R156)
catPrP	(R164), (H187, R185), (K194), (R228), (R229), (K220)
sheepPrP	(R164), (K185), (R228), (H187, K194, H155, R156, R136, R220)
mousePrP[N174T]	(R164), (R229, R230), (K220), (K185, H187, K194) (R156), (H140, R151, R136, R148)
humanPrP[Q212P]-M129	(R164), (K185), (R228)
humanPrP-M129	(R164, H177, K185, H187, K194, H155, R156, R220, R208, R151, R148, H140), (R228)
rabbitPrP[S173N]-NMR	(R163, R135, K184, H186), (R227), (K193, R155), (R147), (H139, R135, R150)
rabbitPrP[I214V]	(R163, R135), (R227), (H186, K184), (K193), (R147), (H139, R135, R150)
rabbitPrP[S170N]-X-ray	(R164, R136, H187, K185, K194), (R228), (H140, R208), (H140), (R148), (R156)
rabbitPrP[S174N]-X-ray	(R164, R136, K185, - H187, R156, K194, R136), (R228), (H140), (R148), (R156)
rabbitPrP[S170N,S174N]-X-ray	(R164, R136, K185, H187, K194, R156, R136), (R228), (H140), (R151), (R148), (R156)
mousePrP at 37 °C	(R164), (R229, R230), (K220), (K185, H187, K194, R136), (R148), (R151)
mousePrP[V166A]	(R164, K185), (R136, K220, R229, R208), (R230), (K194), (R156), (R148), (R151)
mousePrP[D167S] at 20 °C	(R164, K185, H187, K194), (R229, R230), (R156), (R151, R136), (R148), (K220)
mousePrP[D167S,N173K]	(R164), (H177), (R229 R230), (K185, H187), (K194), (R136, R156), (R156), (R148),(R151),(K220)
mousePrP[Y169G]	(R164, K185, K194), (R230), (R229), (K220), (R136), (R156), (R151), (R148)
mousePrP[Y169A]	(R164, K185, K194), (R230), (R229), (K220), (R136, R151), (R156), (R148)
mousePrP[S170N]	(R164), (R230), (R229), (K185), (K194, H187, R136), (R156), (K220), (R208), (R151), (R148)
mousePrP[S170N,N174T]	(R164), (R229), (R230), (K194, K185), (R156), (R136), (R148), (R151), (R208), (K220)
mousePrP[F175A]	(R164, K185), (R229), (R230), (K194, H187, R136, R156, R151), (R148), (H140)
mousePrP[Y225A,Y226A]	(R164, K185), (R229, R230), (R156), (K194, R136), (R148, R151), (R208), (K220)
mousePrP[Y169A,Y225A,Y226A]	(R164), (R229), (R230), (K185, K194), (R136, K220), (R148), (R151)
elkPrP	(R164), (R228), (K185), (HHT121), (K194), (R156), (R136, R220), (R148), (R151)
pigPrP	(R164), (R228), (K185, H187, K194, R156, R136), (K220), (R148, R151)
BankVolePrP	(R164), (R229), (K185), (K194), (R156), (R148), (R151, R136), (H140), (K220)
TammarWallabyPrP	(R164), (R227), (K185, H187), (K194), (R156), (R148), (R136, H140, R148, R151, R156)
rabbitPrP-NMR	(R163, R227), (K184, HHT124)-(K193, R155, R135, R150, H139, R147), (R147)
rabbitPrP-X-ray	(R164, HHT126, K185)-(K194, H187, R136), (R228), (R156), (R148), (R151), (H140)
horsePrP	(R164), (R228,K220,R136),(K194,H187,K185),(R156),(R148, R151),(H177),(K204),(R208),(HHT119)

Table 5:  $\pi$ - $\pi$ -stacking and  $\pi$ -cations for each of the 33 PrPs:

Species	$\pi$ - $\pi$ -stacking	$\pi$ -cations
mousePrP	F175-Y218, Y162-Y128, H187-F198	F141-R208
humanPrP		R164-Y169
bovinePrP		Y128-R164
SyrianHamsterPrP	Y169-F175-Y218	
dogPrP		
catPrP		Y150-R156
sheepPrP	F141-Y150, Y169-F175-Y218	
mousePrP[N174T]	F141-Y150, Y169-F175-Y218	Y128-R164
humanPrP[Q212P]-M129		R228-H237
humanPrP-M129		
rabbitPrP[S173N]-NMR		L124-Y127
rabbitPrP[I214V]-NMR	H139-Y149	Y148-R155
rabbitPrP[S170N]-X-ray	Y169-F175	F141-R208
rabbitPrP[S174N]-X-ray	Y169-F175	F141-R208
rabbitPrP[S170N,S174N]-X-ray	Y169-F175	F141-R208
mousePrP - at 37 °C		R164-Y169
mousePrP[V166A]	Y169-F175	
mousePrP[D167S] - at 20 °C	F175-Y218	
mousePrP[D167S,N173K]	F175-Y218, H187-F198	
mousePrP[Y169G]	F141-Y150, F175-Y218, Y225-Y226	Y128-R164
mousePrP[Y169A]	W145-Y149, H187-F198	
mousePrP[S170N]	Y225-Y226	
mousePrP[S170N,N174T]		R164-Y169
mousePrP[F175A]	Y163-Y218	F141-R208
mousePrP[Y225A,Y226A]	Y169-F175-Y218	
mousePrP[Y169A,Y225A,Y226A] - at 20 °C		
elkPrP	Y169-F175-Y218	
pigPrP		
BankVolePrP	Y169-F175-Y218	R164-Y169
TammarWallabyPrP		R156-F198
rabbitPrP-NMR	F140-Y149	L124-Y127
rabbitPrP-X-ray	Y169-F175	
horsePrP		R156-F198